

Supercontinuum generation in hydrogenated amorphous silicon waveguide

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Abstract: We report supercontinuum (SC) generation centered on the telecommunication C-band (1550nm) in CMOS compatible hydrogenated amorphous silicon waveguide. A broadening of 200nm is obtained in a 1cm long waveguide using as pump 1.6ps pulses with on chip peak power of 4,5W and 45W.

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1. Introduction

Since the first report of supercontinuum (SC) generation in borosilicate glass by Alfano and Shapiro in 1970, many others demonstrations have been successfully reported as for example in single mode fibers [1], photonic crystal fibers [2] and more recently in integrated guiding structures [3].

For the generation of a SC in integrated waveguides both strong nonlinearity and high figure-of-merit (FOM) are required. In crystalline silicon (c-Si) at telecom wavelengths two photon absorption (TPA) clamps the maximum propagating power in the chip and hence reduces dramatically the FOM, which limits the maximum achievable spectral broadening. This problem is overcome by using hydrogenated amorphous silicon (a-Si:H) as nonlinear medium [4].

The high FOM of a-Si:H wires [5] allowed the demonstration of many nonlinear functions such as parametric amplification frequency [4], all-optical signal processing [6], and more recently low power continuous-wave frequency conversion [7]. Here, we present the first demonstration of SC generation in a-Si:H waveguides using telecommunication 1.6-picosecond-long pump pulses.

2. Experiment and results

The a-Si:H waveguides are fabricated in 220-nm-thick hydrogenated amorphous silicon deposited on top of silicon dioxide layer using a low temperature Plasma Enhanced Chemical Vapor Deposition (PECVD) process. The a-Si:H film was formed by plasma decomposition of silane (SiH₄) gas and combined with Helium (He) for dilution. Waveguides of varying lengths L (1 to 7cm) were fabricated using 193 nm optical lithography and dry etching. The cross-section of these waveguides is 500×220 nm². To allow an effective injection and extraction of the light from the photonic nanowires, broadband grating couplers are added at the input and the output. Thanks to a cutback measurement at low power we estimate propagation losses for TE polarization of about 3dB/cm.

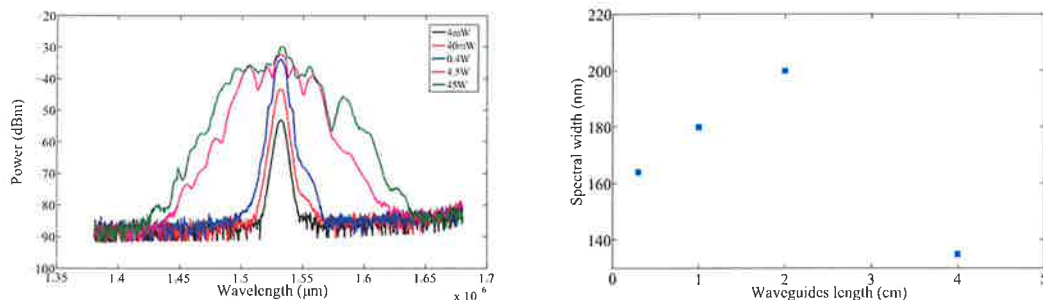


Figure 1: (a) Transmission spectra as function of on chip peak pump power at the exit of a 1cm long waveguide. (b) Spectral width of the supercontinuum as a function of the waveguide length ($P_{\text{peak}}=45\text{W}$).

For the supercontinuum generation experiment, a picosecond pulse train (FWHM ~ 1.6ps, repetition rate = 10 MHz) generated by a tunable optical fiber cavity is used as the pump. The coupling loss between the cleaved fiber and the waveguides through the in/out grating coupler is estimated to be 6-7dB. The output beam is sent into an optical

spectrum analyzer (OSA). The polarization is aligned with the TE mode of the waveguide with the help of polarization controllers.

Figure 1(a) shows the dependence of the output spectrum as a function of on chip peak pump power at pump center wavelength of $\lambda_{\text{pump}}=1530\text{nm}$. At low pump power, the spectral width varies only slightly. However, above a threshold pump power the spectral width increases sharply until at the highest power, 45W, the spectral width has increased to reach a value close to 200nm. The measured spectral broadening is restricted by the bandpass output grating coupler (200nm at -20dB), and that the spectral broadening on chip may be much larger. We note that the spectral width is determined from the intersection of the signal with the noise level of our OSA.

To understand the nonlinear processes involved in this widening, we investigate the evolution of the output spectrum as a function of the pump power. At an input peak power of 4.5W, the pump has undergone self-phase modulation and two modulation instability sidebands have appeared detuned by approximately 65nm around the pump wavelength of 1530nm (fig. 1a). At the highest peak power of 45W, cascaded four wave mixing (FWM) has broadened and flattened the spectrum further.

Using the intensity of the converted signal in a continuous wave FWM experiment [8], we estimated the nonlinear coefficient ($n_2=1.7\times 10^{-17}\text{m}^2/\text{W}$) and dispersion coefficient ($\beta_2=2\text{ps}^2/\text{m}$) of the waveguide. This implies a nonlinear length $L_{\text{NL}}=83\mu\text{m}$ and soliton fission length $L_{\text{fission}}=1.03\text{cm}$ [2]. These values confirm that the FWM is the main process in our supercontinuum generation.

Finally we investigated the influence of the waveguide length (L) on supercontinuum generation using L=3mm, 1cm, 2cm and 4cm long waveguides (fig. 1b). We observe that the spectral broadening increases with waveguide length until L=4cm, where the linear losses are very high which clamps output spectrum widening.

3. Conclusion

We have demonstrated SC generation in 1cm long CMOS compatible hydrogenated amorphous silicon waveguides. We measured broadening of 200nm is a lower bound on the width of the supercontinuum as it is limited by the bandpass of the output grating coupler. In the future experiments with no output grating coupler will be performed to reveal the maximum width of the SC.

4. References

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ABSTRACT BODY:

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9:00 PM - 9:15 PM NM3A.6. Infrared supercontinuum generation in soft-glass photonic crystal fibers
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9:15 PM - 9:30 PM NM3A.7. Supercontinuum generation in hydrogenated amorphous silicon
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